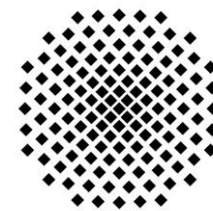




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DEGLI STUDI
FIRENZE

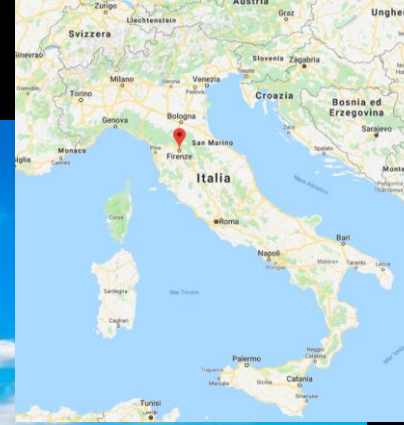
MOLECULAR SPIN QUBITS STUDIED BY EPR AND THZ SPECTROSCOPY

Lorenzo Tesi



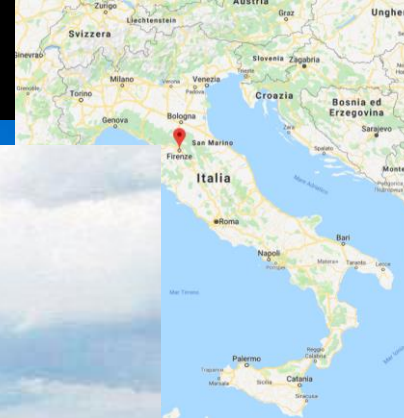
Universität
Stuttgart

19-06-2019

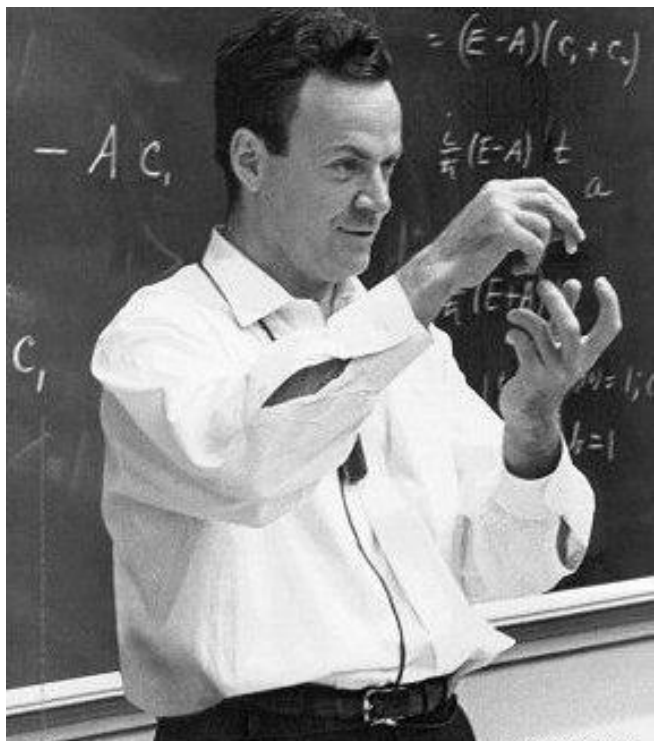


My Home





WHAT IS A QUBIT?



Richard Feynman, 1982

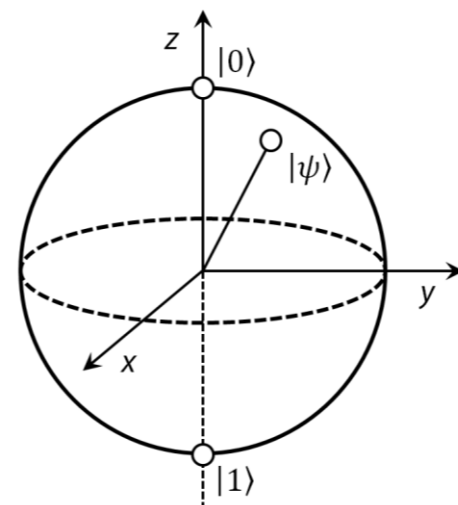
'If you want to make a simulation of nature you'd better make it quantum-mechanical'



Classical bit

$$|\psi\rangle = |0\rangle \text{ or } |1\rangle$$

Two possible states



Quantum bit

$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

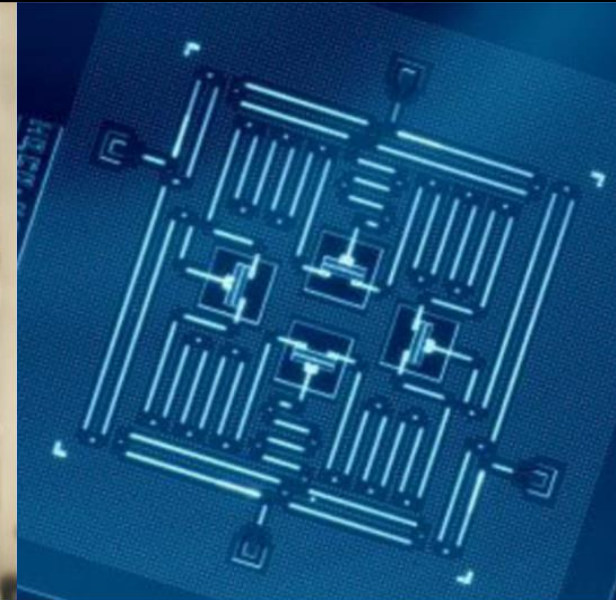
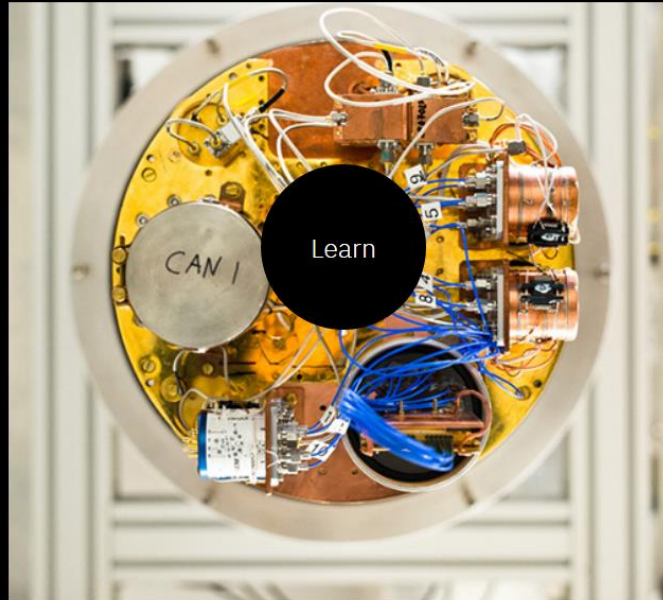
Superposition state

Faster computation and data searching capabilities;
Quantum systems can be simulated.

PROTOTYPE OF QUANTUM COMPUTER



Quantum
Computing



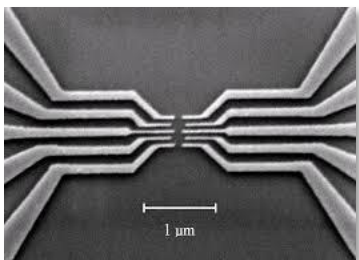
Meet IBM Q

IBM Q is an industry-first initiative to build commercially available universal quantum computers for business and science. While technologies like AI can find patterns buried in vast amounts of existing data, quantum computers will deliver solutions to important problems where patterns cannot be found and the number of possibilities that you need to explore to get to the answer are too enormous ever to be processed by classical computers. We invite you join us in exploring what might be possible with this new and vastly different approach to computing.

www.research.ibm.com

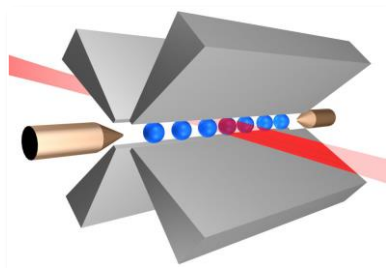
But also Google, Nasa, USA Energy Department, Alibaba, etc...

WHAT IS A QUBIT?



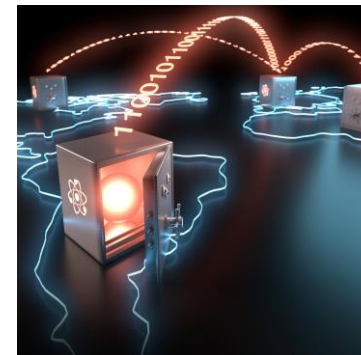
Quantum Dots

Nature **453**, 1043-1049



Ionic Traps

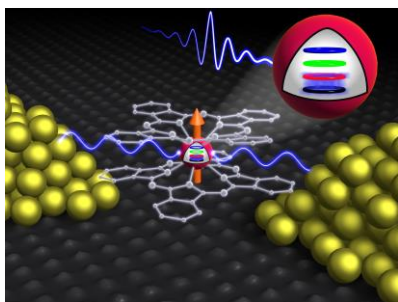
Nature **453**, 1008-1015



Photons

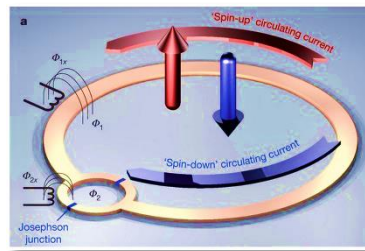
Nature **409**, 46-52

POTENTIAL QUBITS



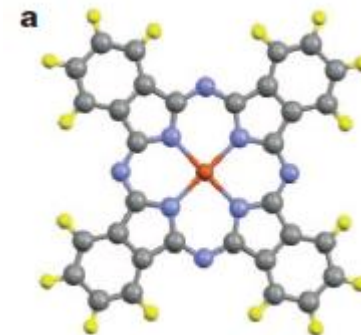
Nuclear Spins

Nature **496**, 334-338



Superconducting Circuits

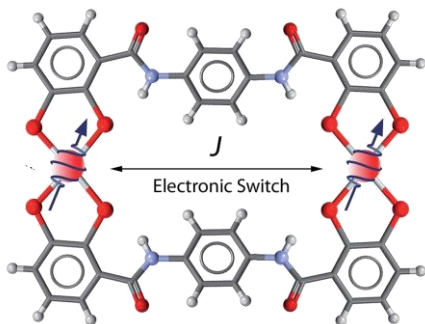
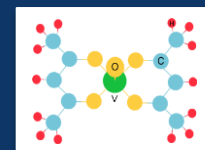
Nature **453**, 1031-1042



Electronic Spins

Nature **503**, 504-508

ADVANTAGES AND DRAWBACKS



- ✓ Easily Interacting
- ✓ Scalable
- ✓ Easy quantum error correction

BUT **short** coherence times

$$\begin{array}{rcc}
 \Delta m_l = -1 & |\Phi_-\rangle & \xrightarrow{\alpha} \text{---} \text{---} \\
 \Delta m_l = 0 & |\Phi\rangle & \xrightarrow{\alpha} \text{---} \text{---} \\
 \Delta m_l = +1 & |\Phi_+\rangle & \xrightarrow{\alpha} \text{---} \text{---} \\
 m_l & & \begin{array}{ccc} \frac{5}{2} & \frac{3}{2} & \frac{1}{2} \end{array}
 \end{array}$$

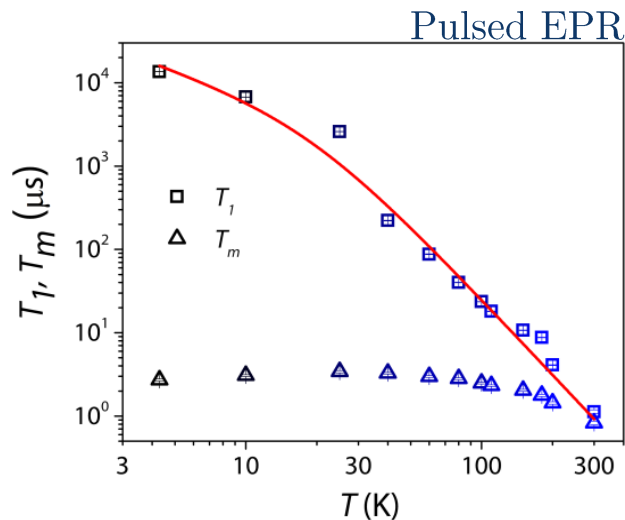
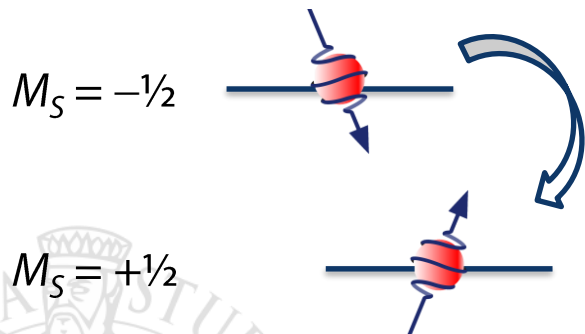
Spin-lattice
Relaxation Time T_1



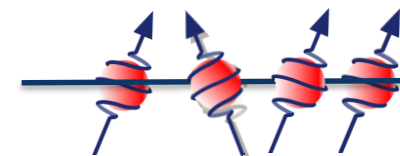
LIMITING FACTOR

Phase
Memory Time T_m
(or Coherence Time)

Classical
Information
Memory



Quantum
Information
Memory

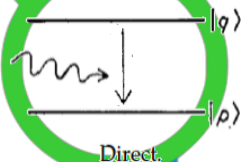


$$|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$$

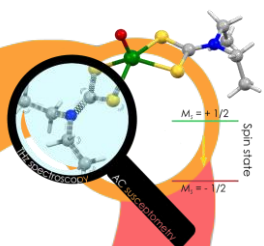
We need to identify the **ingredients** to lengthen the relaxation times!



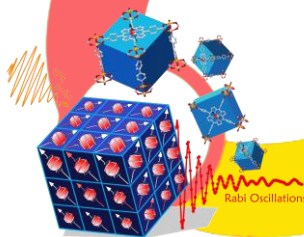
Part I
Spin-lattice Relaxation
Theoretical inputs



Part III
New evidences
from Terahertz spectroscopy

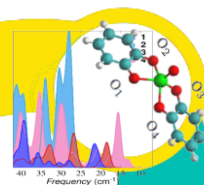


Part II
An example of
Molecular Spin Qubits

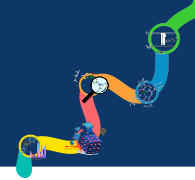


Part V
Developing a
theoretical framework

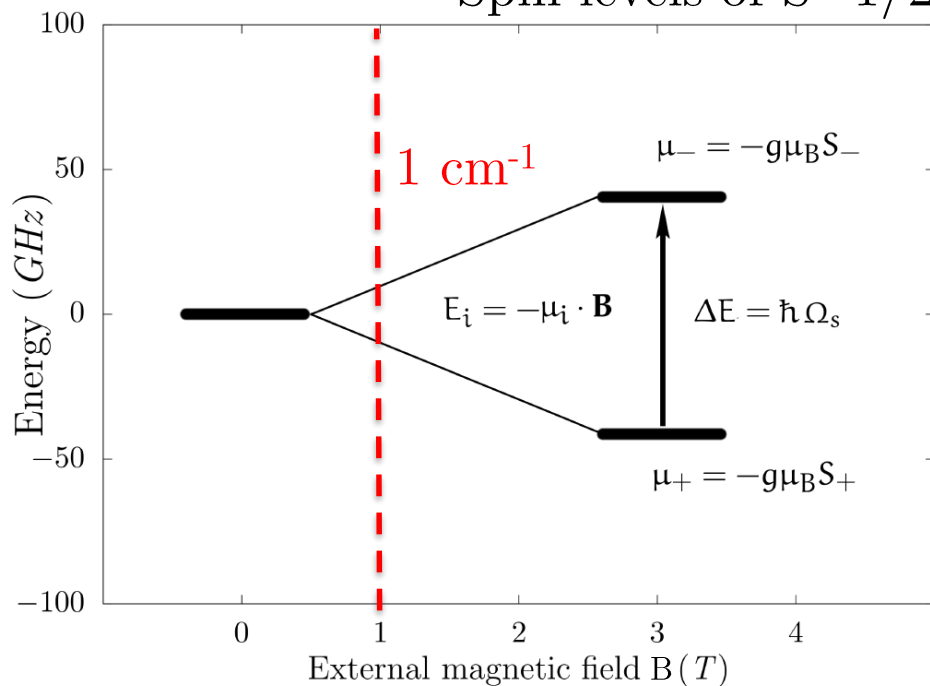
Part IV
Quantum MOF



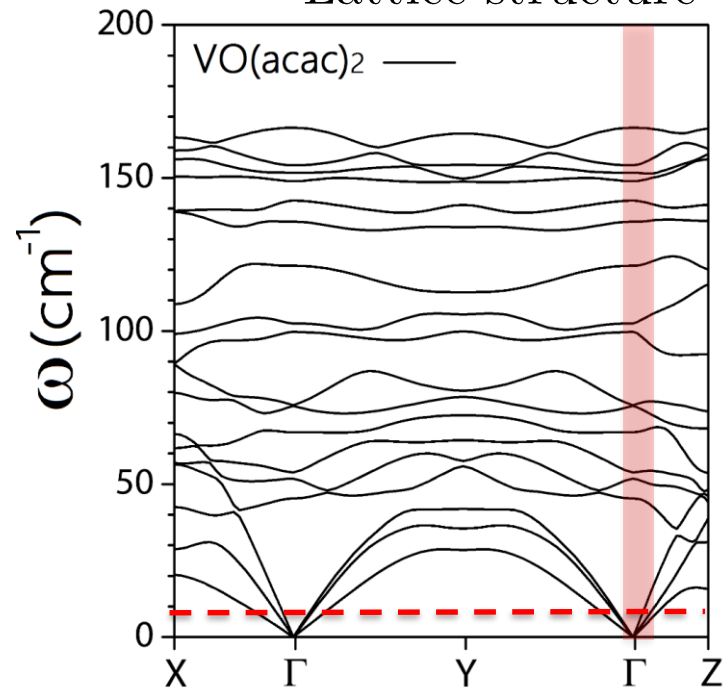
CONCLUSIONS



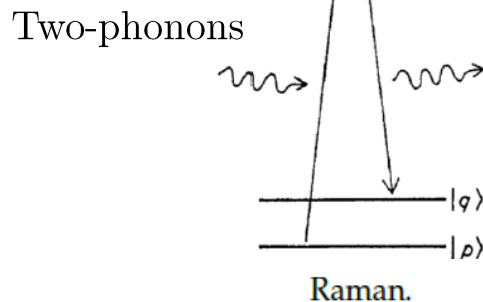
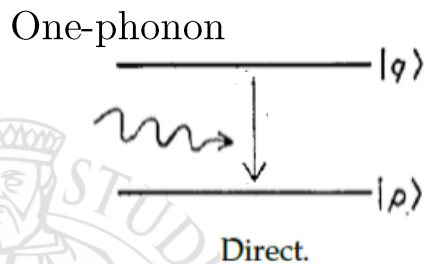
Spin levels of $S=1/2$



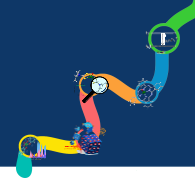
Lattice structure



Main relaxation processes

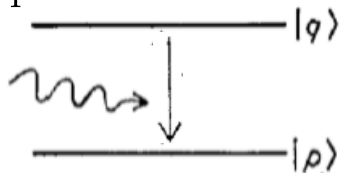


The efficiency of spin-lattice relaxation is given by the **spin-orbit coupling**.



Main relaxation processes

One-phonon



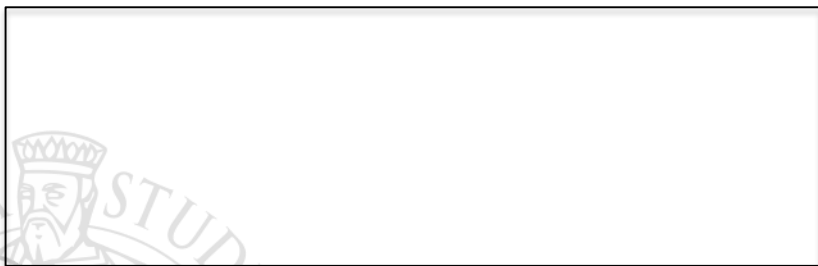
Direct.

$$\frac{1}{T_1^D} = \frac{3}{2\pi} \left(\frac{\delta_{pq}}{\hbar} \right)^3 \frac{\coth(\delta_{pq}/2kT)}{\rho \hbar v^5} W_D^2$$

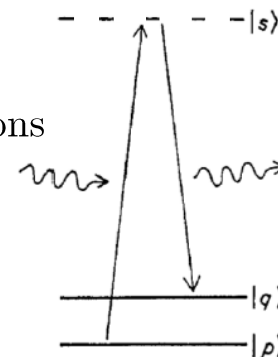
Temperature dependence Direct

$$\frac{1}{T_1} = aT$$

Magnetic field dependence of T_1



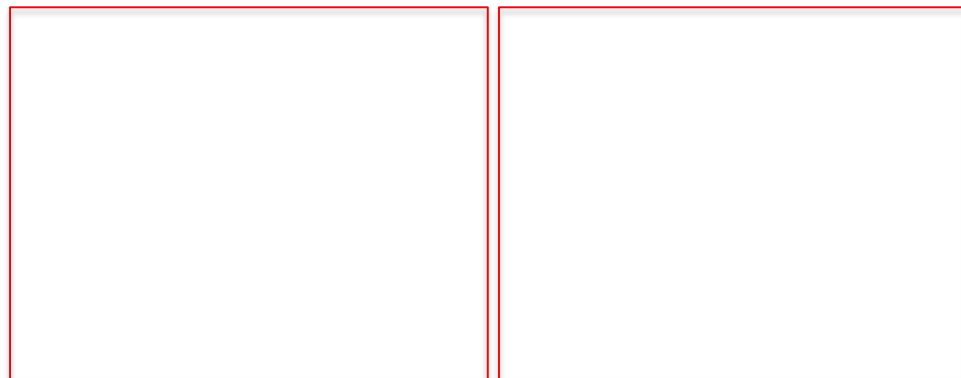
Two-phonons

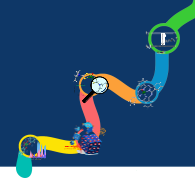


Raman.

$$\frac{1}{T_1} \sim 8! \left(\frac{kT}{\hbar} \right)^9 \frac{W_R^2}{\Delta^2} \left(\frac{2V}{4\pi^2 M v^5} \right)^2$$

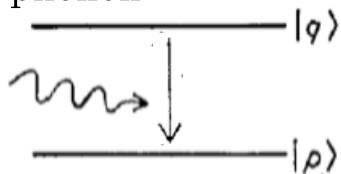
Temperature dependence Raman





Main relaxation processes

One-phonon



Direct.

$$\frac{1}{T_1^D} = \frac{3}{2\pi} \left(\frac{\delta_{pq}}{\hbar} \right)^3 \frac{\coth(\delta_{pq}/2kT)}{\rho \hbar v^5} W_D^2$$

Temperature dependence Direct

$$\frac{1}{T_1} = aT$$

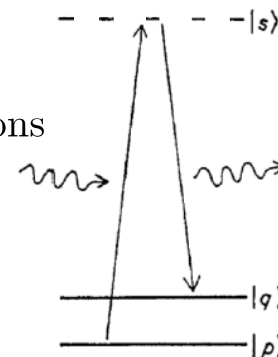
Magnetic field dependence of T_1

$$\frac{1}{T_1} = cB^4 + d \frac{1 + eB^2}{1 + fB^2}$$

direct

Brons-van
Vleck

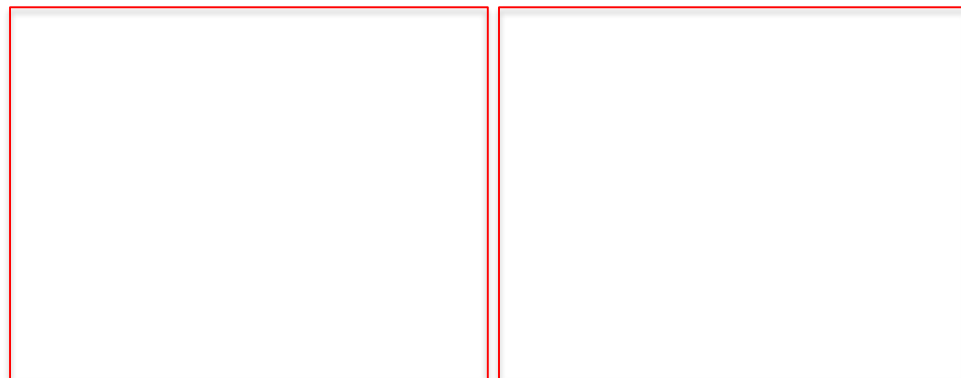
Two-phonons

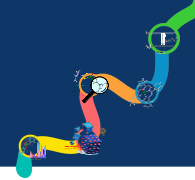


Raman.

$$\frac{1}{T_1} \sim 8! \left(\frac{kT}{\hbar} \right)^9 \frac{W_R^2}{\Delta^2} \left(\frac{2V}{4\pi^2 M v^5} \right)^2$$

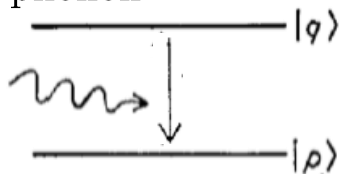
Temperature dependence Raman





Main relaxation processes

One-phonon



Direct.

$$\frac{1}{T_1^D} = \frac{3}{2\pi} \left(\frac{\delta_{pq}}{\hbar} \right)^3 \frac{\coth(\delta_{pq}/2kT)}{\rho \hbar v^5} W_D^2$$

Temperature dependence Direct

$$\frac{1}{T_1} = aT$$

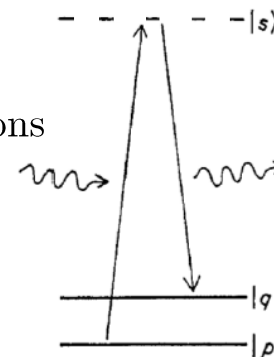
Magnetic field dependence of T_1

$$\frac{1}{T_1} = cB^4 + d \frac{1 + eB^2}{1 + fB^2}$$

direct

Brons-van
Vleck

Two-phonons



Raman.

$$\frac{1}{T_1} \sim 8! \left(\frac{kT}{\hbar} \right)^9 \frac{W_R^2}{\Delta^2} \left(\frac{2V}{4\pi^2 M v^5} \right)^2$$

Temperature dependence Raman

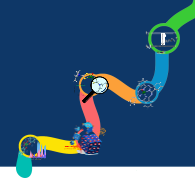
Debye model

$$\frac{1}{T_1} = a_{Ram} \left(\frac{T}{\theta_D} \right)^9 I_8 \left(\frac{\theta_D}{T} \right)$$

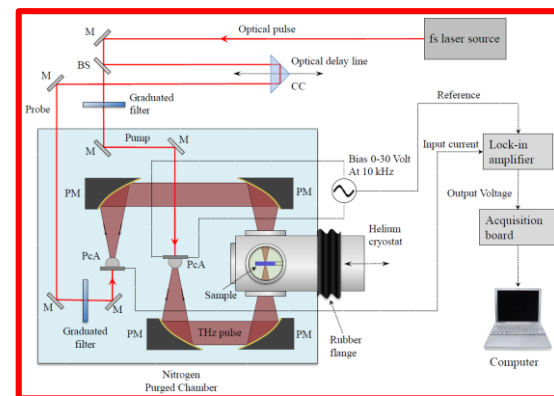
$$I_8 \left(\frac{\theta_D}{T} \right) = \int_0^{\theta_D} \frac{x^8 \exp x}{\exp(x-1)^2} dx$$

Local modes model

$$\frac{1}{T_1} = a_{loc} \frac{\exp(\hbar\omega)}{\left(\exp\left(\frac{\hbar\omega}{kT}\right) - 1 \right)^2}$$



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CONTINUOUS-WAVE
EPR



Static
Magnetic
Properties

AC
SUSCEPTOMETRY



Magnetization
Relaxation
Time τ

PULSED EPR

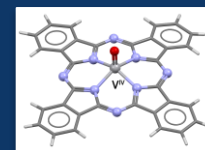


Spin-Lattice
Relaxation Time T_1 ;
Phase Memory Time T_m

THZ SPECTROSCOPY



Low-energy
vibrational spectrum

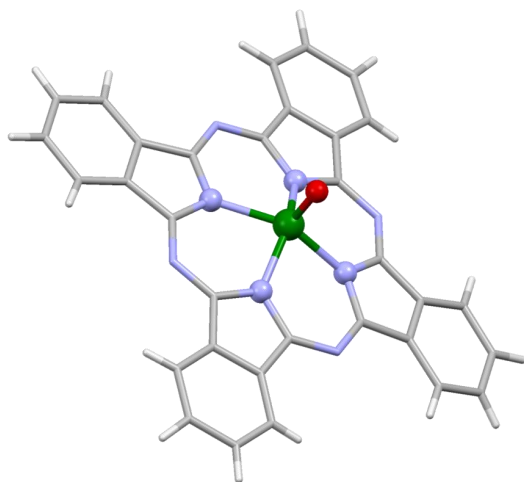

 M. Atzori, L. Tesi *et al.*, *JACS*, 2016

$$S = 1/2 \quad I = 7/2$$

$$\mathcal{H} = \mu_B \mathbf{g} \cdot \hat{\mathbf{S}} \cdot \mathbf{B} + \hat{\mathbf{S}} \cdot \mathbf{A} \cdot \hat{\mathbf{I}}$$

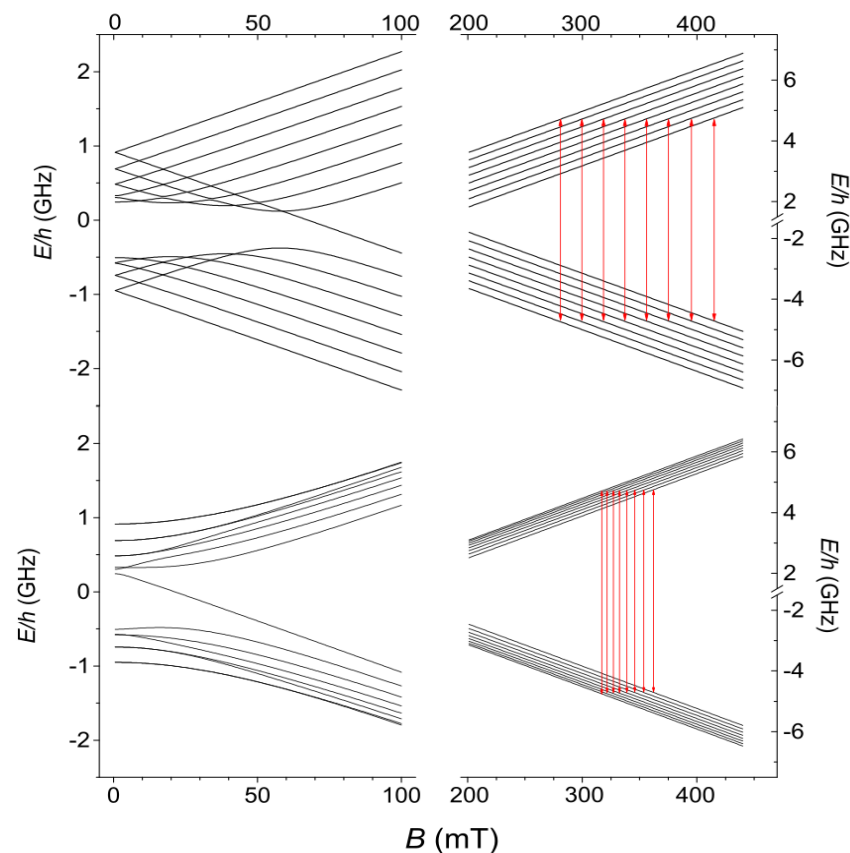
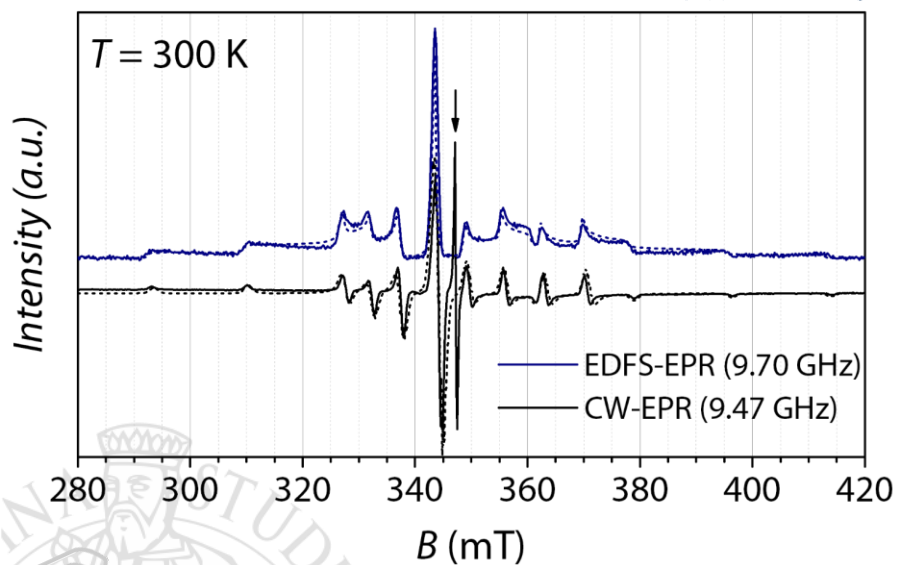
$g_x = g_y$	g_z	$A_x = A_y$	A_z
1.987(1)	1.966(1)	167.9(1) MHz	476.7(2) MHz

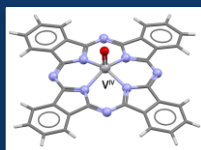
Spin levels composition



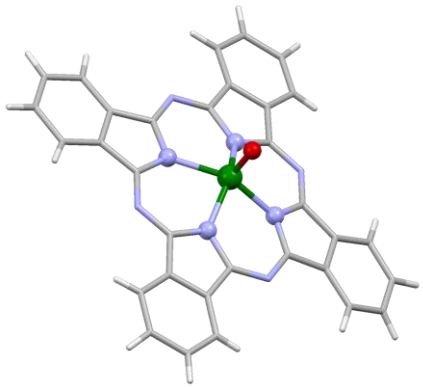
VOPc 0.1% in TiOPc

cw-EPR





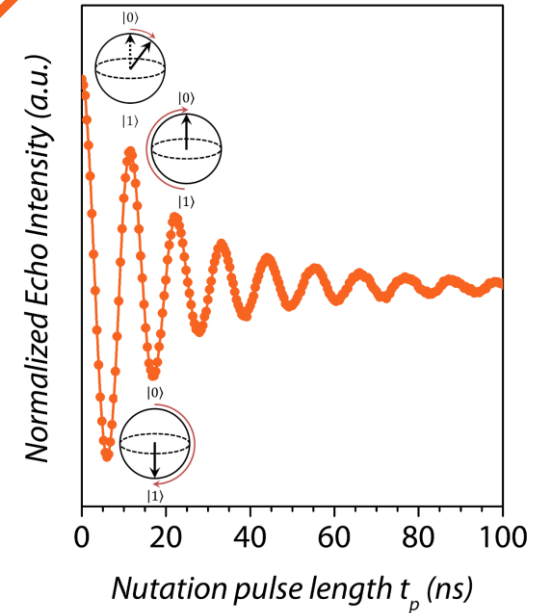
M. Atzori, L. Tesi *et al.*, *JACS*, 2016



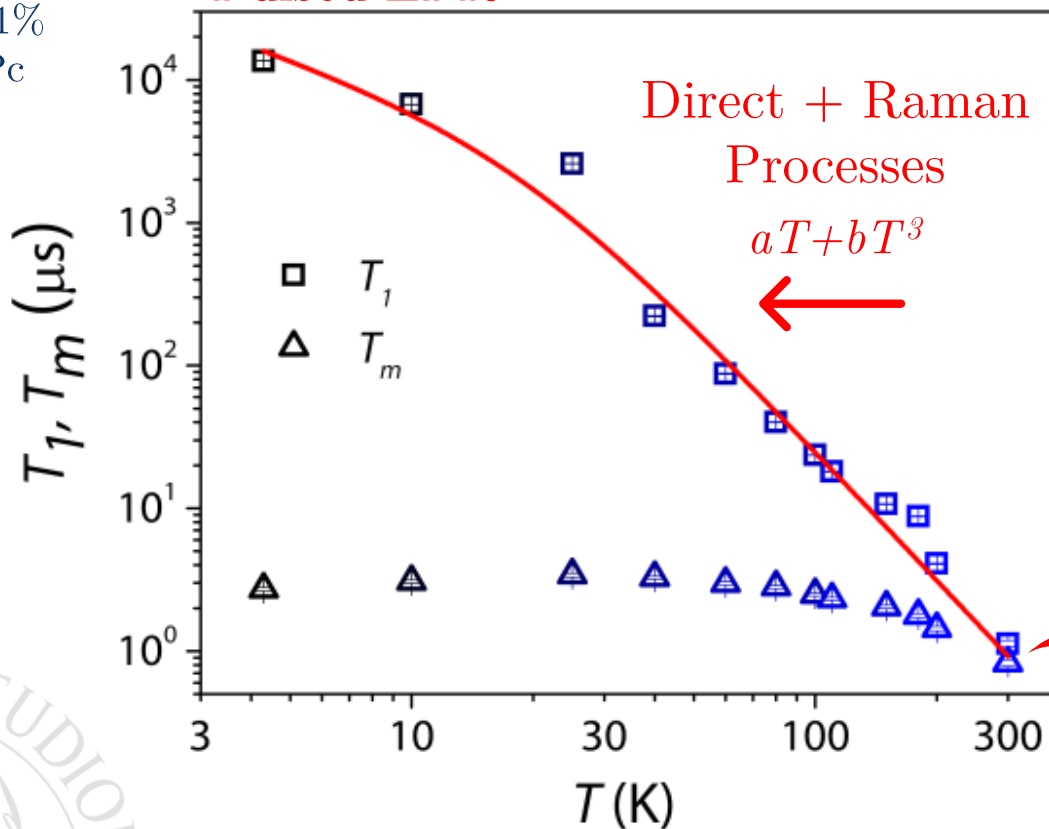
VOPc 0.1%
in TiOPc

Quantum
manipulation →

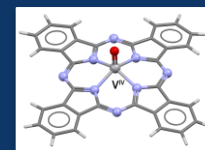
Rabi oscillations



Pulsed EPR



0.83 μs
at 300 K

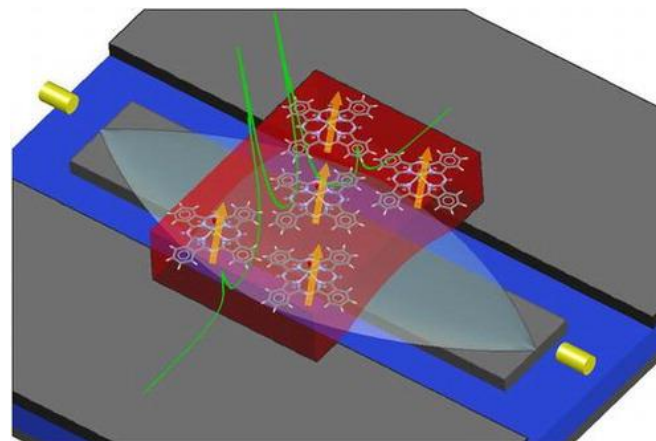


Room Temperature
coherence of
VOPc drew attention
for other applications...

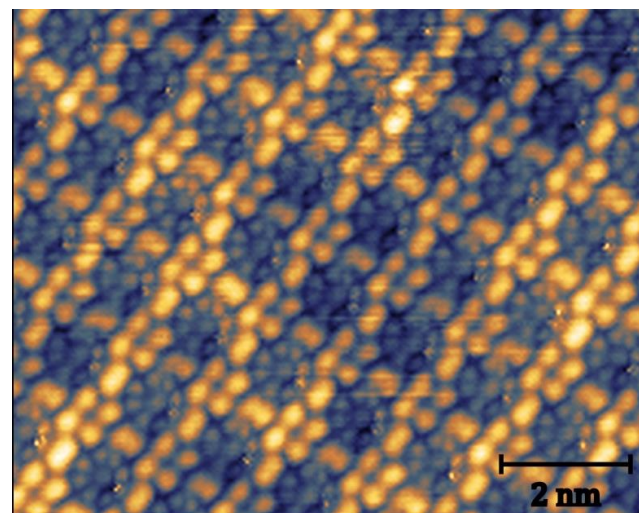
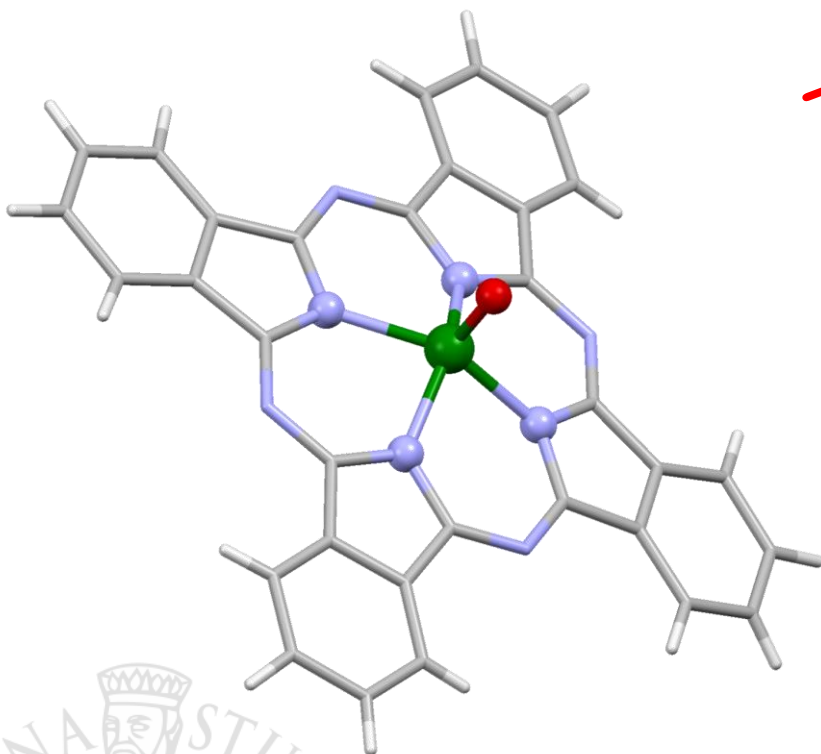


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MODENA E REGGIO EMILIA

Superconducting planar resonator



C. Bonizzoni et al. *Sci. Reports* 2017



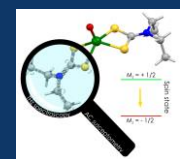
I. Cimatti et al. *Nano. Horizons* 2019

L. Malavolti et al. *Nano Letters* 2018

M. Atzori, L. Tesi et al., *JACS*, 2016



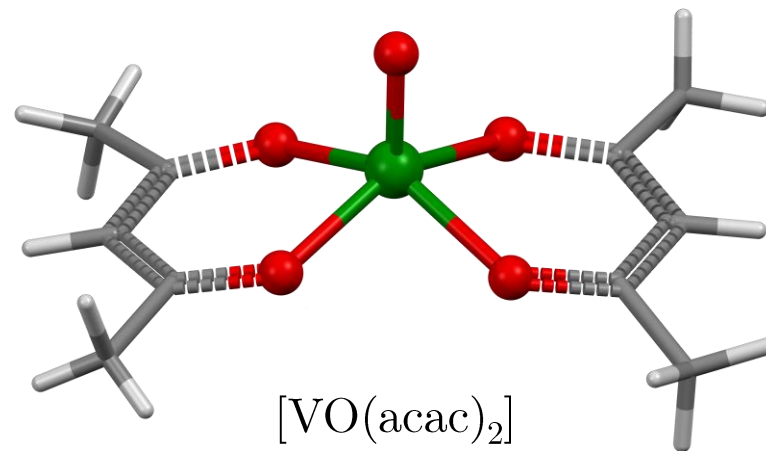
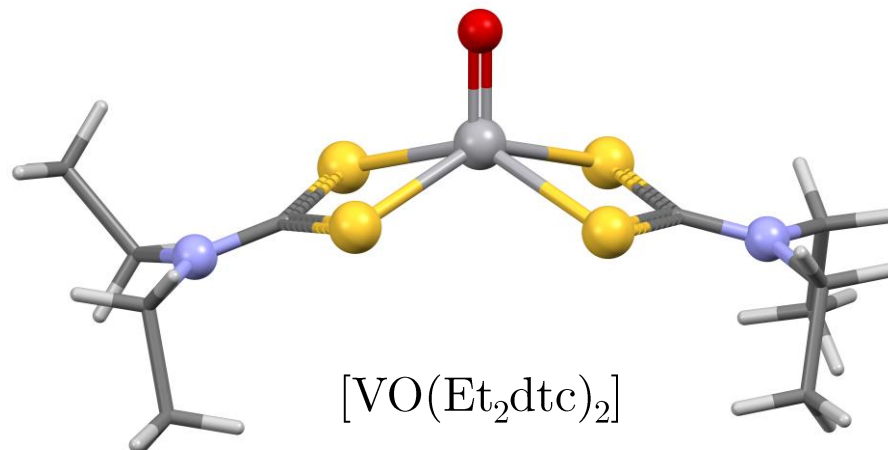
Max Planck Institute



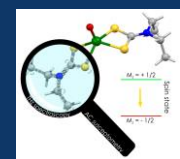
AC susceptometry and
THz spectroscopy
of the pure compounds



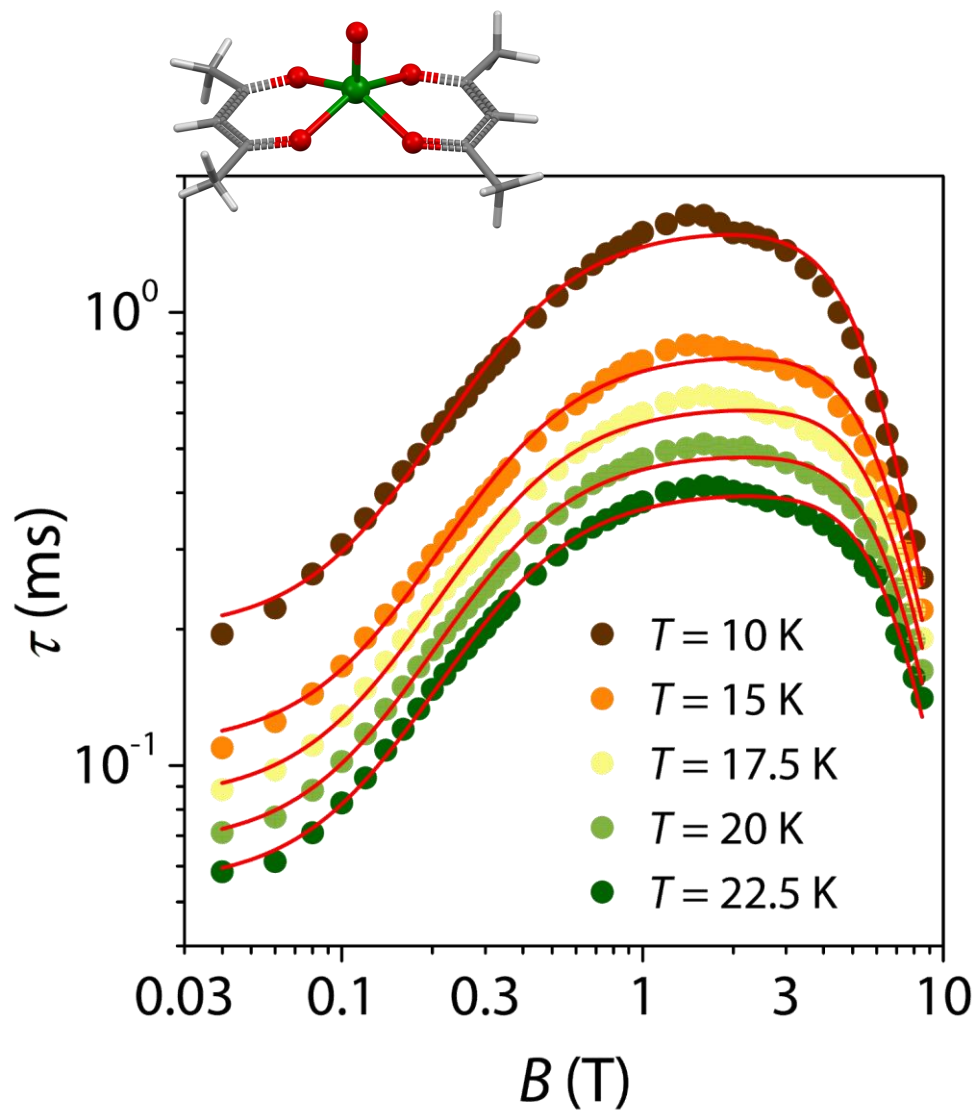
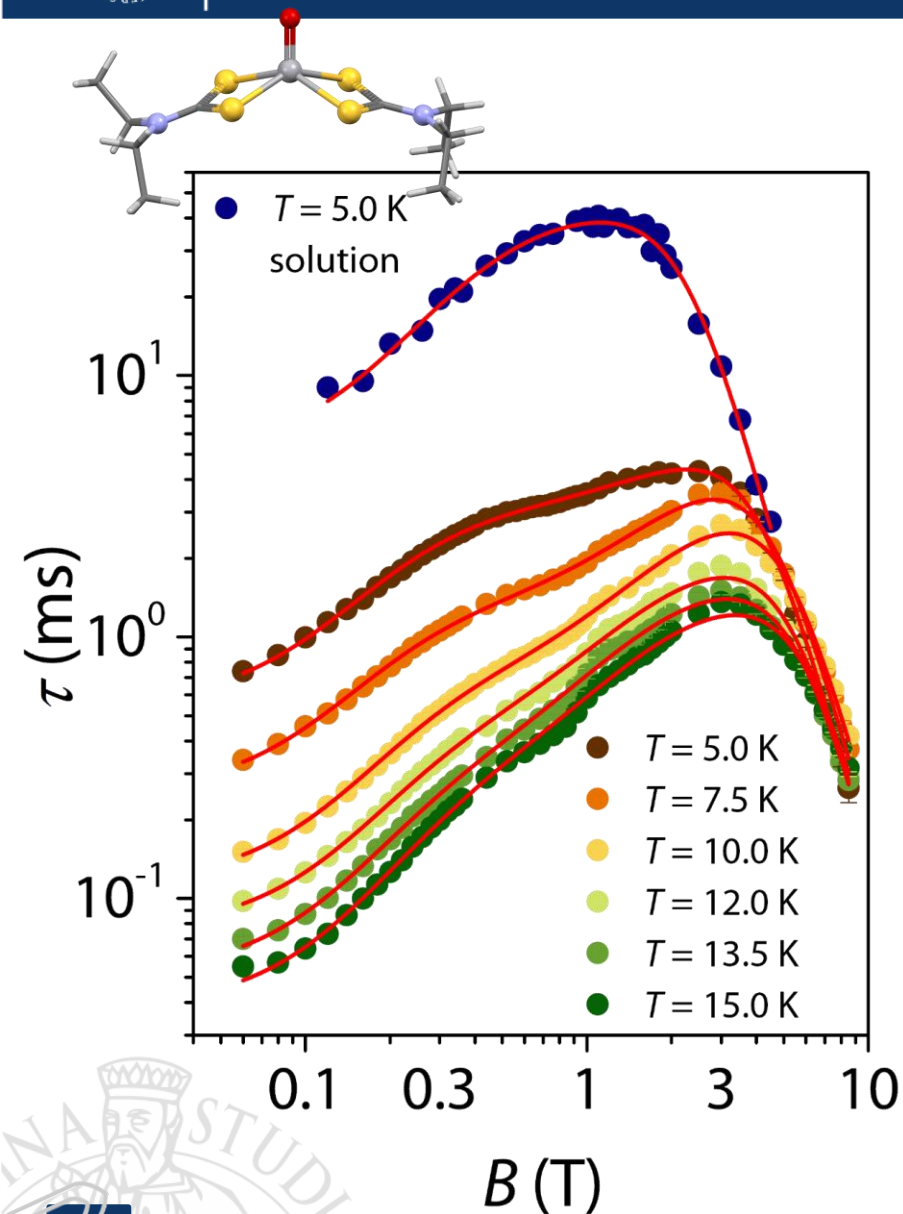
Evidence of the role of **low energy phonons** in spin dynamics

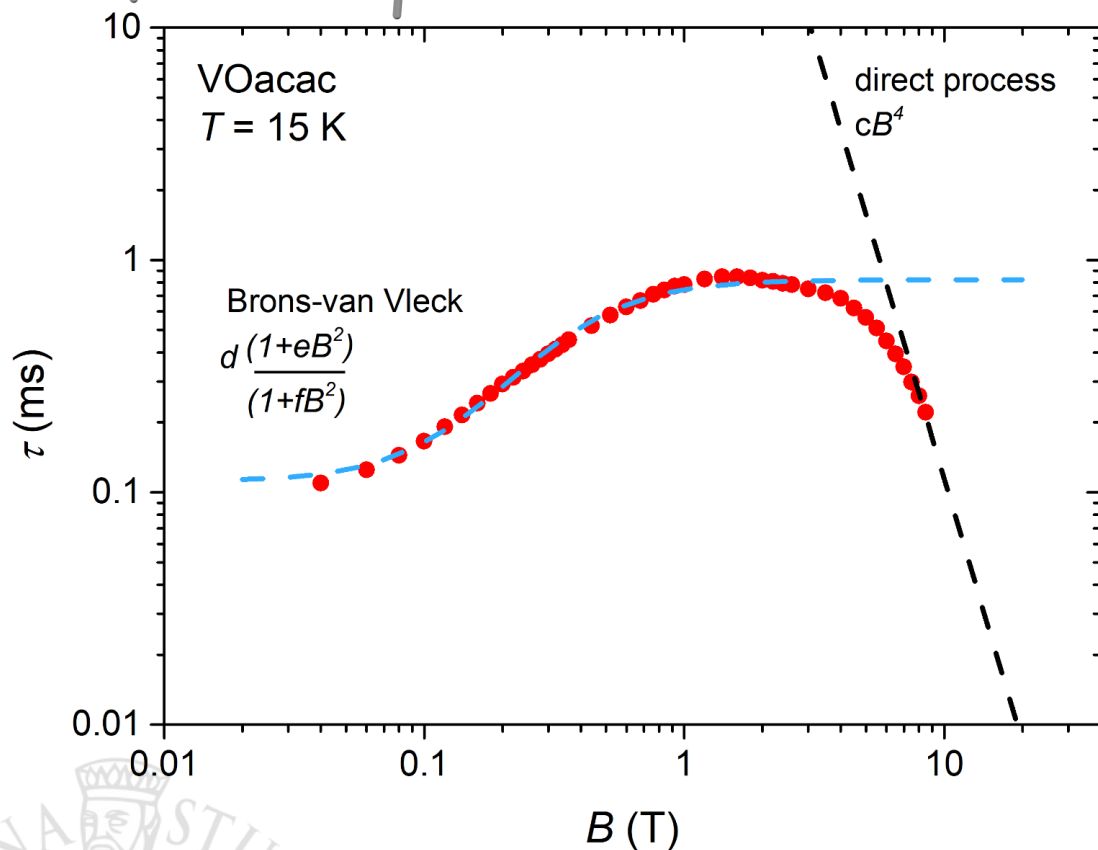
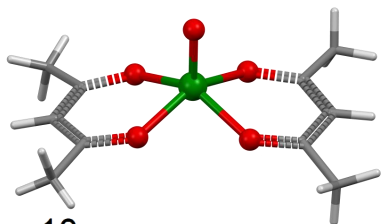
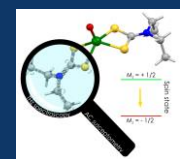


(Et_2dtc = Ethyldithiocarbamate)
(acac = acetylacetonate)



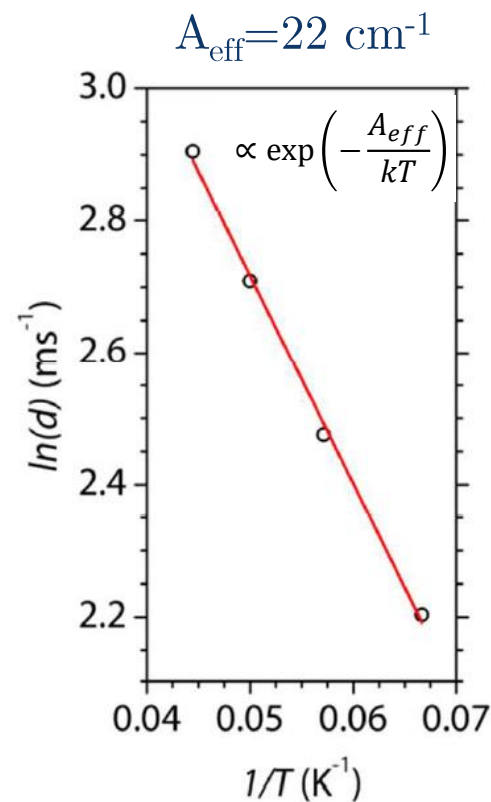
AC susceptometry – B dependence of τ

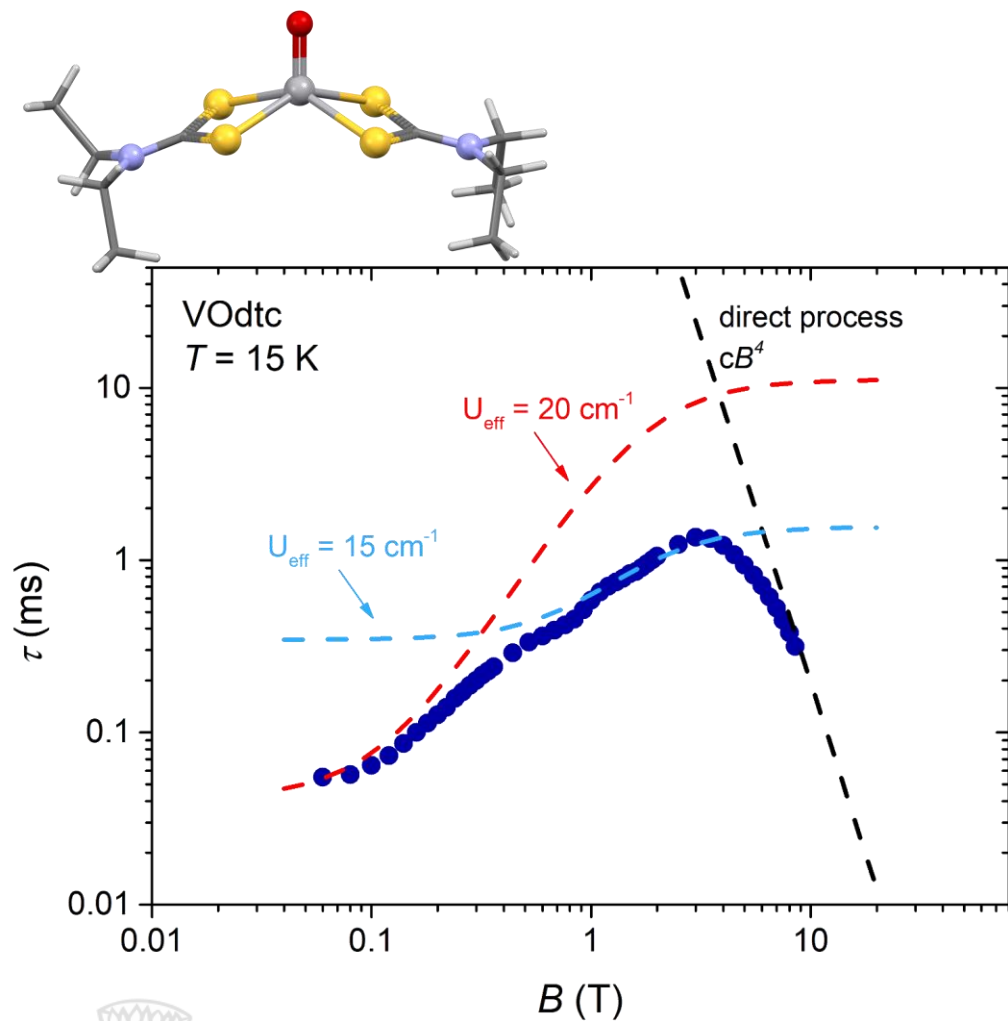
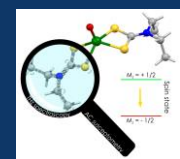




Direct + Brons-van Vleck model

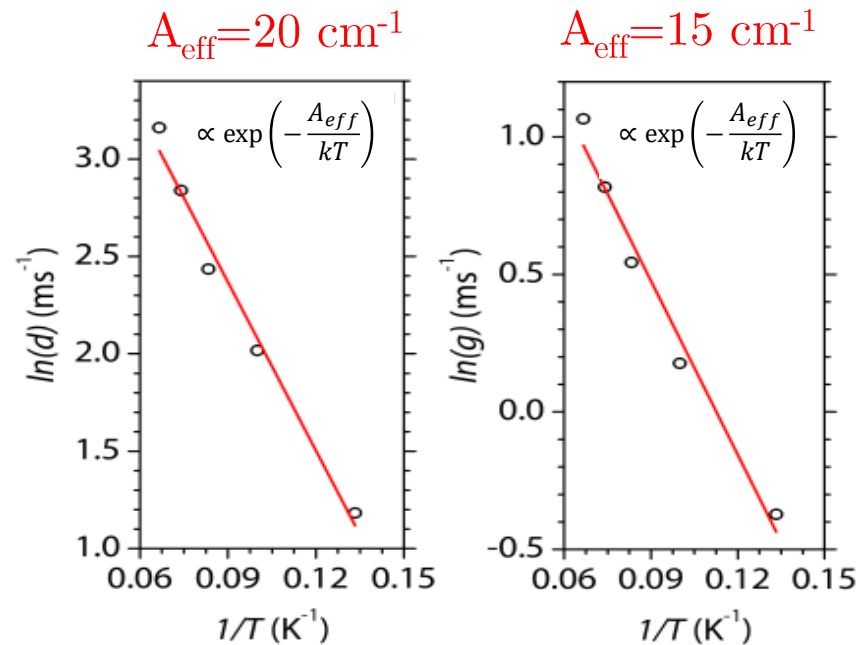
$$\tau^{-1} = cB^4 + d \frac{1 + eB^2}{1 + fB^2}$$

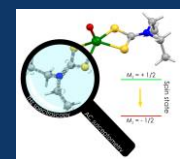




Direct + Brons-van Vleck model

$$\tau^{-1} = cB^4 + d \frac{1 + eB^2}{1 + fB^2} + g \frac{1 + eB^2}{1 + hB^2}$$

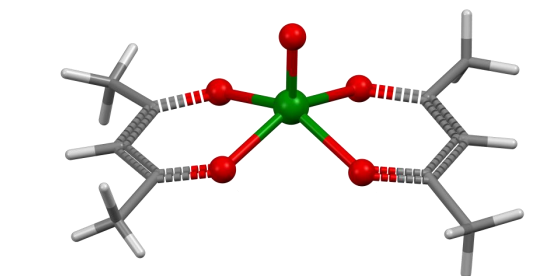




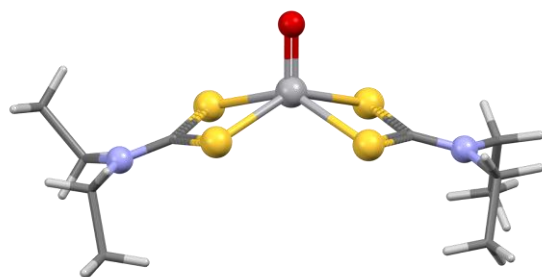
Correlation between the activation energy and phonon frequency:

From: A. Lunghi *et al.*, *Nature Comm.* 8, 14620 (2017)

$$A_{eff} = \frac{h\omega_v}{2}$$



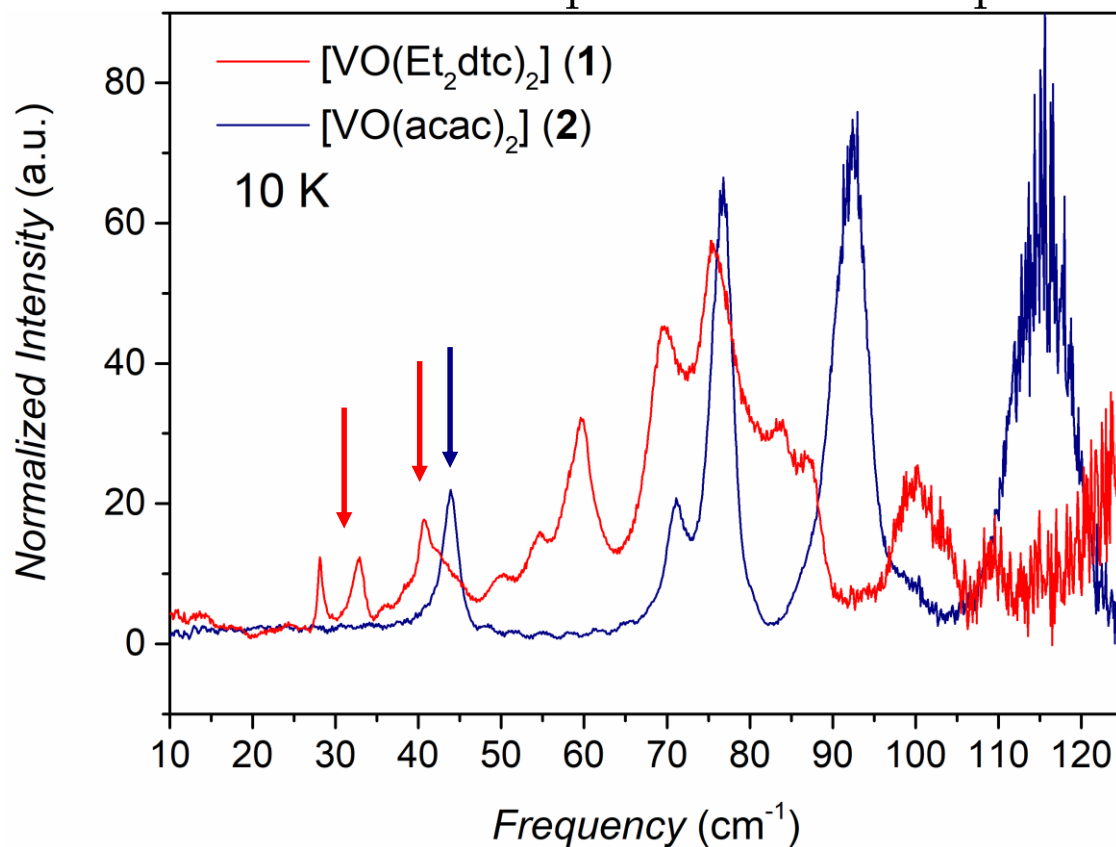
$$A_{eff} = 22(2) \text{ cm}^{-1}$$

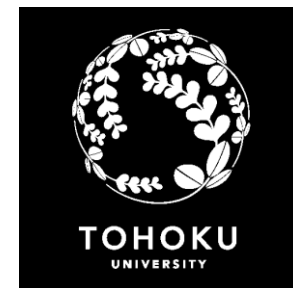
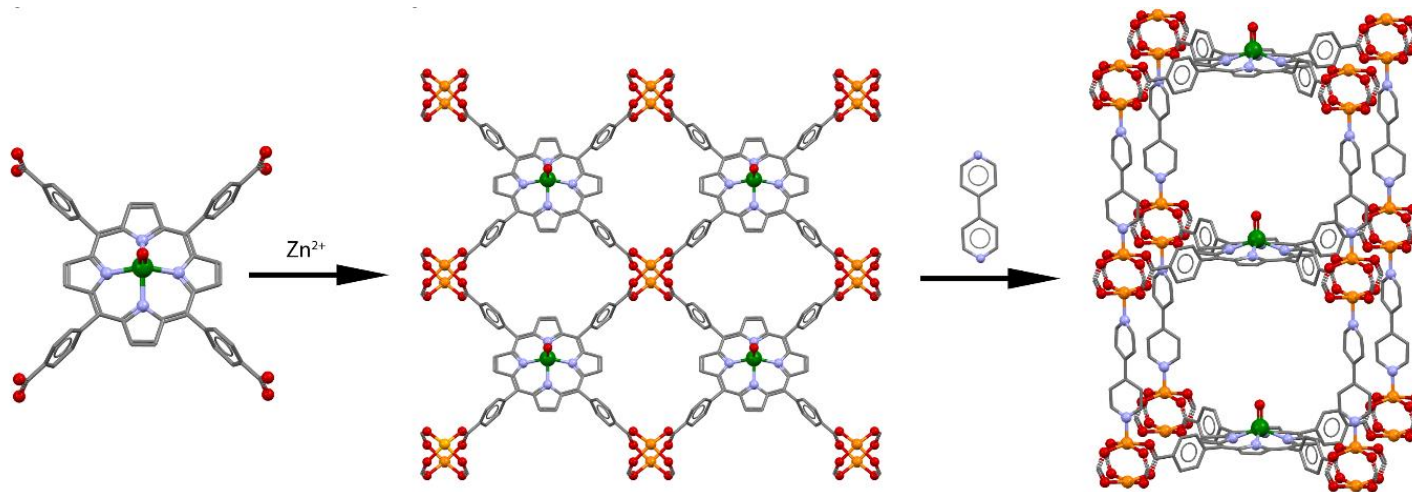
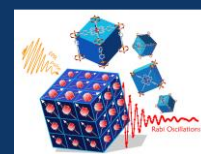


$$A_{eff} = 15(2) \text{ cm}^{-1}$$

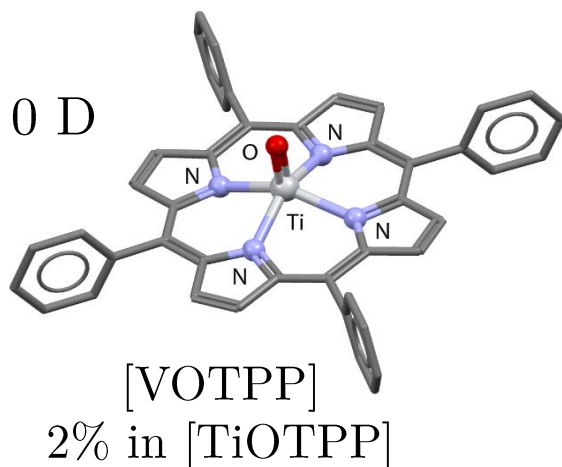
$$A_{eff} = 20(2) \text{ cm}^{-1}$$

experimental THz spectra

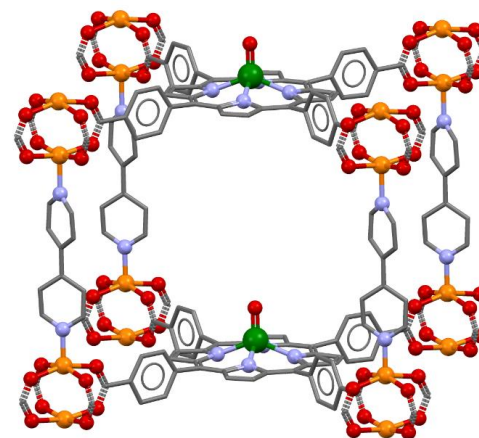




In collaboration with
Prof. M. Yamashita
and T. Yamabayashi



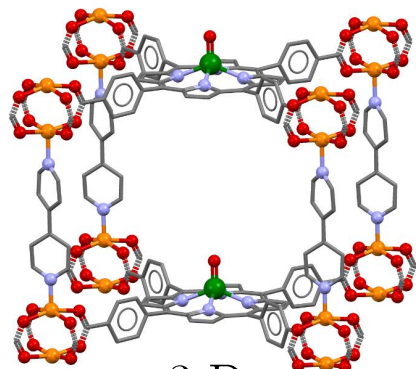
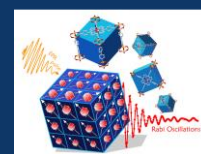
3 D



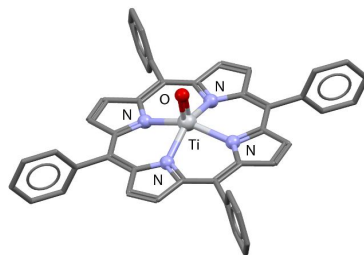
[VO(TCPP-Zn₂-bpy)]
5% in [TiO(TCPP-Zn₂-bpy)]

(VOTPP = tetraphenylporphyrinate)
(TCPP = tetracarboxylphenylporphyrinate)
(bpy = 4,4'-bipyridyl)

T. Yamabayashi, M. Atzori, L. Tesi *et al.*, *JACS*, 2018



3 D



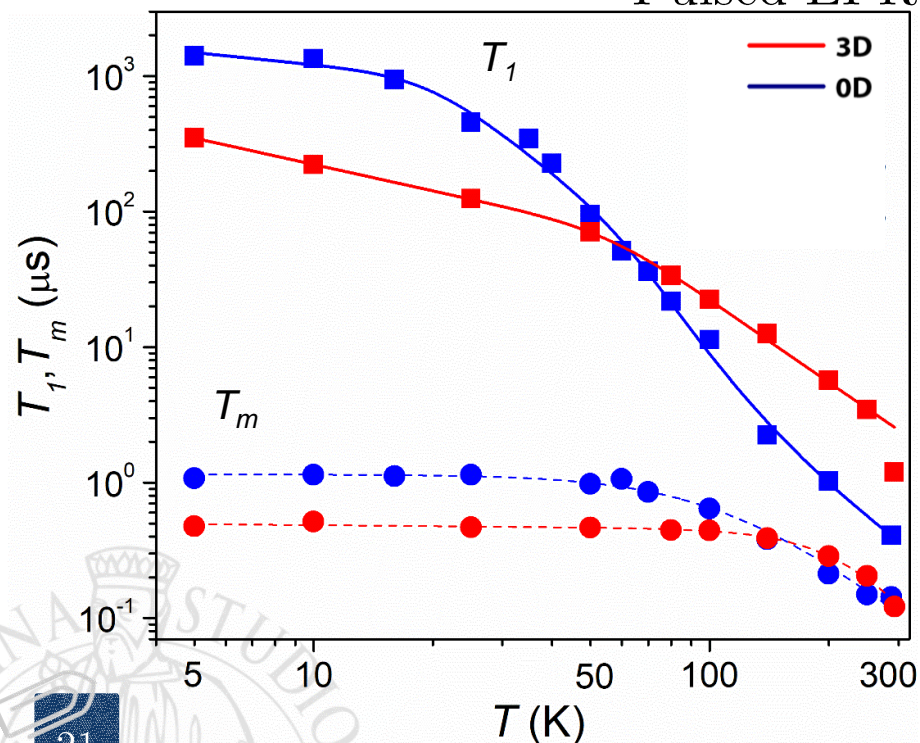
0 D

$$T_1^{-1} = a_{dir} T^x + a_{loc} \frac{e^{(\hbar\omega/k_B T)}}{(e^{(\hbar\omega/k_B T)} - 1)^2}$$

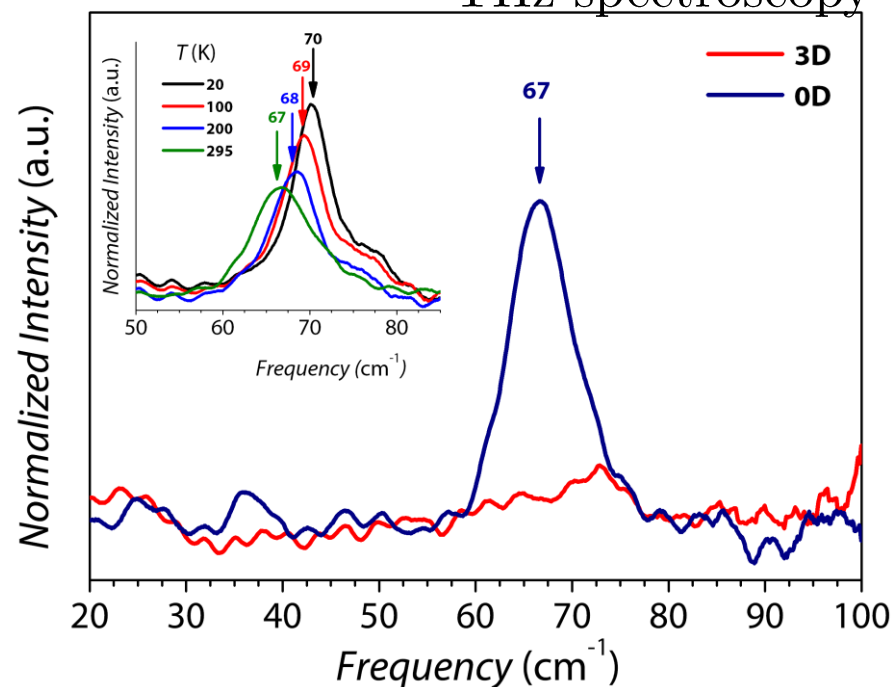
$$3D \quad \hbar\omega = 184(2) \text{ cm}^{-1}$$

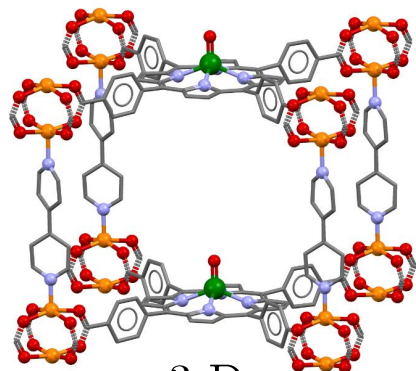
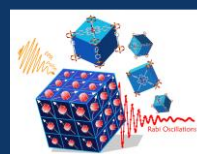
$$0D \quad \hbar\omega = 67(0) \text{ cm}^{-1} + \hbar\omega = 303(35) \text{ cm}^{-1}$$

Pulsed EPR

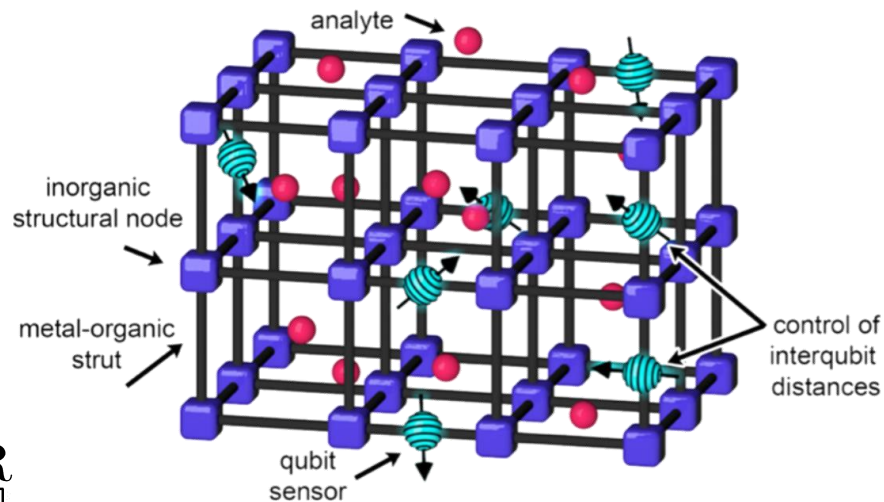


THz spectroscopy

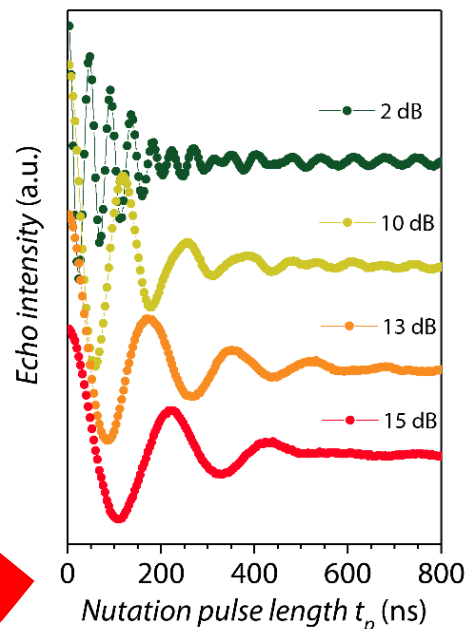
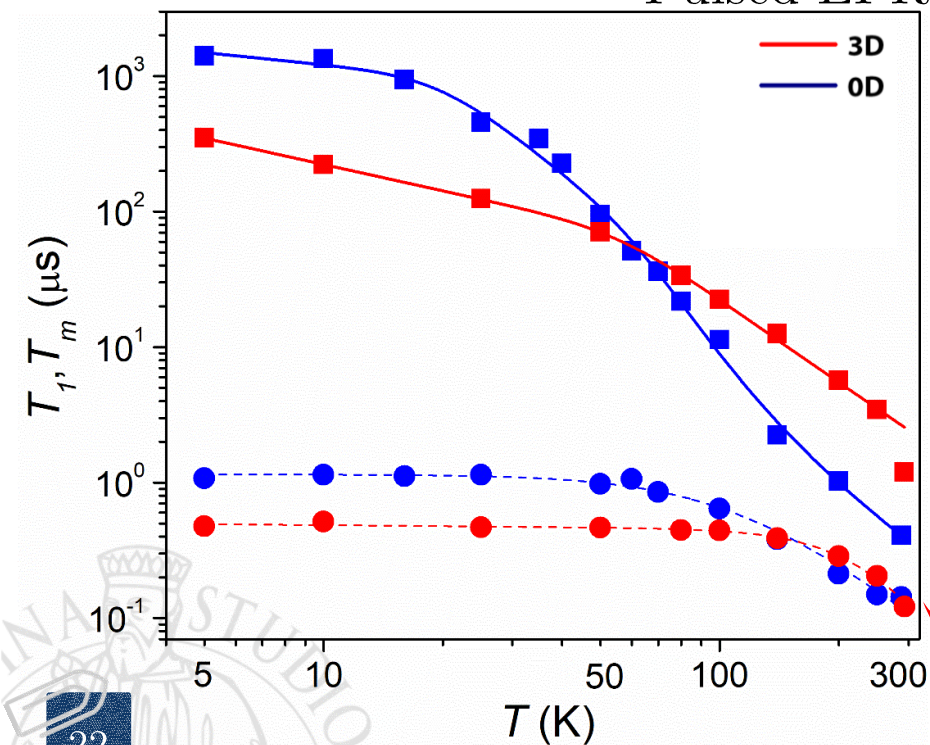




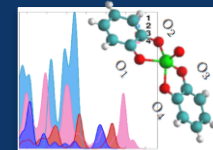
3 D



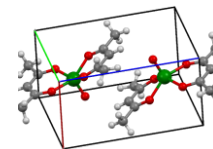
Pulsed EPR



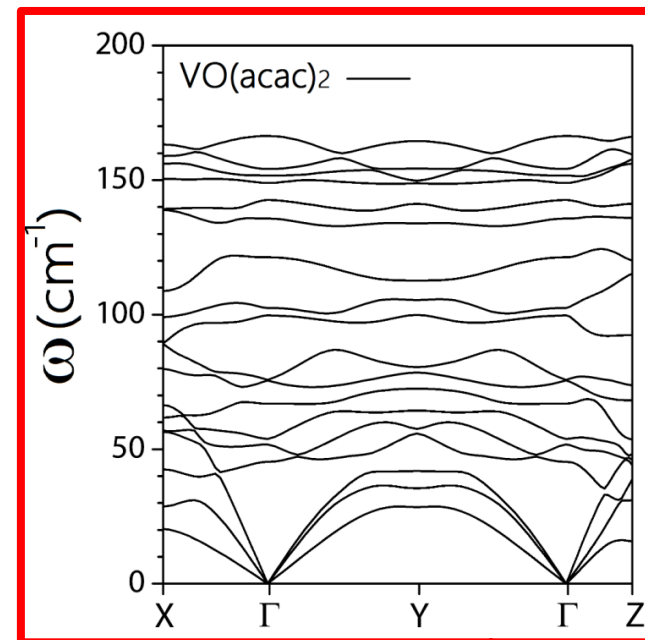
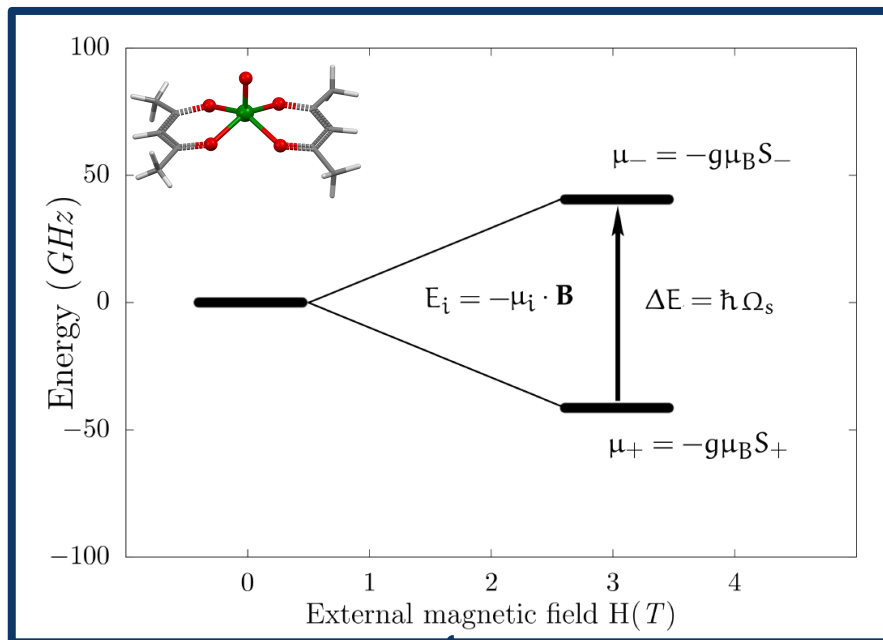
Rabi oscillations



Spin system



Lattice system



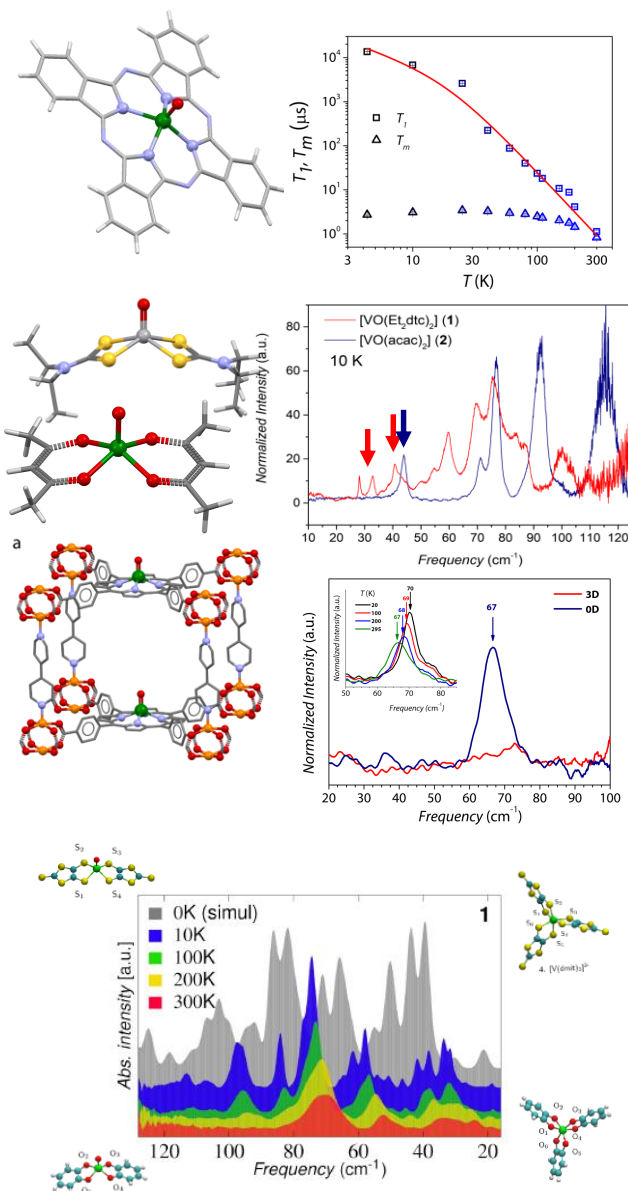
A. Albino, S. Benci, L. Tesi *et al.*, *ArXiv*, 2019
A. Lunghi, S. Sanvito, *ArXiv*, 2019

**Spin-phonon
coupling**

by *ab initio* methods

CONCLUSIONS

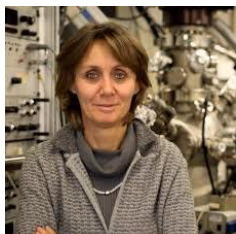
- Vanadyl phthalocyanine shows long coherence times also at room temperature. This, together with the possibility to easily deposit it on surface makes it an interesting MSQ system;
- We have evidenced the role of low energy vibrational modes in spin dynamics by combining the results of ac susceptibility (τ vs B) and pulsed EPR (T_1 vs T) measurements with THz spectroscopy;
- Ab initio calculations can be used to simulate the spin-phonon coupling, this is an important step toward a better understand of the spin-lattice relaxation mechanisms. It is important to check the quality of the simulation in the THz range.





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The Florence group



Prof. R. Sessoli Prof. L. Sorace Dr. M. Atzori F. Santanni

THz spectroscopy



Prof. R. Torre

S. Benci



UNIVERSITÀ
DEGLI STUDI
DI TORINO

Pulsed EPR

AND
YOU



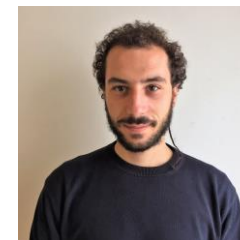
Trinity
College
Dublin

Theoretical calculations



Prof. M. Chiesa Dott. E. Morra

FOR
THE



Dr. A. Lunghi

A. Albino

ATTENTION!